REFORMED HYDRAULIC BULGING IN PROCUREMENT OF STRESS - STRAIN CURVE

ABINASH PANDA, Raajdhani Engineering College, Bhubaneswar DR.PRADYUT KUMAR SWAIN, Aryan Institute of Engineering and Technology, Bhubaneswar DR. KASINATH DASMOHAPATRA, NM Institute of Engineering and Technology, Bhubaneswar Mr. BIBHUTI SABAT, Capital Engineering College, Bhubaneswar

ABSTRACT:

From old style hydraulic bulging of a circular diphagram, another strategy for recording a pressure strain bend is inferred named changed hydraulic bulging on the circle. New strategy permits more exact assurance of film pressure in the misshaping material. Supposition on the mass calculation is at this point not required since adjustment, utilizing circle as a deterrent, decides the sweep of bend required in securing of layer pressure. Performing changed hydraulic bulging upon only one test example, yields three marks of pressure strain bend. Tentatively got focuses have various strains, covering various regions; little, medium and huge disfigurements, inside a solitary trial. Utilization of the little transducer coordinated in device, empowers obtaining of power and along these lines dependable pressure procurement upon the known (round) math. Similarly, the acquired pressure strain bend can be utilized in a limited component model of the changed hydraulic bulging on the circle that is currently a work in progress

KEYWORDS: Modified hydraulic bulging, stress - strain curve, hydroforming

1 INTRODUCTION

From the experimental work on hydraulic bulging, performed on various sheet materials [1, 2] and from analytical considerations based on plasticity theory [3], the method has been constantly developing through experimental simplifications and automatization, in order to speed up stress and strain acquisition [4]. Hydraulic bulging has never attained the popularity of standard acquisition methods of true stress-strain curves such as tensile or compressive tests [5]. One of the reasons to this is dependency of the true stress upon geometrical assumptions set on bulge geometry. Many authors have conducted detailed research on geometry of the bulge and thickness distribution [6], some of the putting the emphasis on strain determination [7]. Giving more importance to the calculation of stress, led to the introduction of modified hydraulic bulging on the sphere.

2 CLASSICAL HYDRAULIC BULGING

2.1 IMPORTANCE OF KINEMATICS

Kinematics of hydraulic bulging is very important, since the stress depends on bulge geometry and kinematics. Stress is derived from membrane stress assumption [8]

$$\underline{\sigma_m} + \underline{\sigma_c} = \underline{p}$$

where σ_m = meridional stress, r_m = meridional radius, σ_c = circular stress, r_c = circular radius, p = forming pressure

and *s* = sheet thickness. Assumptions at the pole; $\sigma_m = \sigma_c = k_f$ and $r_m = r_c = r$ give stress as the function of geometry and pressure *p*.

where k_f = flow stress, r = radius at the pole. Radius at the pole can be determined in two ways

It can be calculated from known height of the bulge h and bulging diameter D [3] assuming spherical geometry of the bulge

$$1 a^{2}$$

r = 2 (h + h) (3)

Radius can be calculated by measuring the coordinates of three points near the pole [4]

In classical hydraulic bulging, schematically shown in Figure 1, blank is fixed around its perimeter in a such way, that only its inner part of diameter D is plastically deformed by the pressure p.

Fixation should disable any draw in, in order to provide plane strain $\theta = 0$ around the perimeter of *D*. This pre-

 requisite is theoretically taken for granted, but practically it can be completely accomplished and thus both ways of

pole radius acquisition become erroneous.

Moreover neither exact deforming diameter D cannot be obtained correctly, since edge of the die at the point F has to be rounded with some radius Figure 1, in order to avoid material splitting on the sharp edge. Smaller corner

^{*} Corresponding author: Ivana Lucica 5, HR-10000 Zagreb, CROA-TIA, +385 95 8936 123, marko.skunca@fsb.hr

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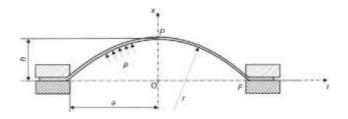


Figure 1: Schematic representation of classical hydraulic bulging

radius results in more accurate determination of *D*, enabling more precise determination of pole radius. Since minimal corner radius is material dependent, standardization of classical hydraulic bulging is not simple. It should also include the standardization of the mentioned radius. Modified hydraulic bulging completely avoids this problem.

The only prerequisites in modified hydraulic bulging are conditions that have to be met in order to have membrane stress state in the deforming sheet;

- Sheet has to be thin regarding deforming radius;
 R > 20s
- 2. Meridional stress gradient has to be small
- 3. Change in the sheet curvature has to be moderate

σт,к =

All this assumptions can be found in where mathematical form [8], while only the first was strictly used in construction of

central part of the tooling, for 1 mm thick sheet as shown in Figure 5.

3 MODIFIED HYDRAULIC BULGING

The only difference between classical and modified hydraulic bulging is in the sphere, positioned above the deforming blank as shown in Figure 2. Deforming pressure

p "pushes" sheet onto the sphere with a $\sigma = \frac{p_K}{c_{,K}}$ force *F* that

as acquired by the force transducer positioned in the axis of symmetry above the sphere.

3.1 STRESS AT POINT K

Going from the axis of symmetry x, K is the point where the sheet separates from the sphere i.e. contact and force

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exertion upon the sphere ceases at point *K*. From the geometry of the deformed blank; radius r_{K} , angle ϑ_{K} shown in Figure 3 and sheet thickness at the point *K*, true meridional stress $\sigma_{m,K}$ can be calculated at point *K* as

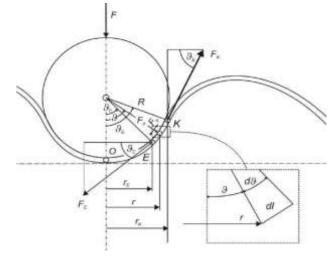


Figure 3: Detailed representation forces and geometry in calculation of the stress at point E



$$F_{\mathcal{K}} = F - \rho r^2 \pi \tag{5}$$

where F = force obtained by transducer, p = deforming pressure, r_{K} = radius at point K (Figure 3), A_{K} = throughthickness surface at point K and θ_{K} = angle between axis of symmetry and normal at point K. Since point K lies on the sphere, radial and meridional curvature are the same, so from equation (1) circular stress $\sigma_{c,K}$ can be calculated

$$SK \qquad R = \sigma_{m,K} \qquad (6)$$

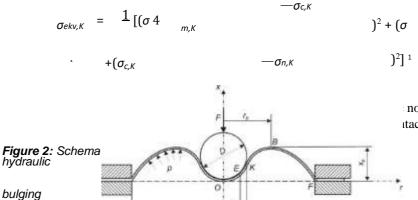
where s_{κ} = thickness at point K, R = radius of the sphere and p_{κ} = contact pressure at point K calculated as

$$p_{K} = \frac{F}{K} p$$
 (7)

where r_{κ} = radius at point *K* (shown in Figure 3). Equivalent von Misses stress at point *K* $\sigma_{ekv,K}$ calculated from meridional and circular stress is equal to

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3.2 STRESS AT POINT E

Point *E* is the point of maximal thinning of the sheet in modified bulging. Therefore, equilibrium of forces in radial direction between points K and E gives,

$$\sum_{r=0}^{\Sigma} F_r = 0$$

$$F_{E} \underline{\cos}(\theta_{E}) \ _F_{tr} \cos(\theta) + F_{\rho} \sin(\theta) + F_{\kappa} \cos(\theta_{\kappa}) = 0$$
(11)

resulting in force at point *E* calculated as

$$F_E =$$

$$\cos(\theta_{E})$$

$$\cos(2\theta_{\rm E})$$
] + 2p_K

$$(\sigma m, K)^{2} + (\sigma m, K)^{2}$$
 (8)

normal stress at point K calculated as mean tact pressure and deforming pressure p;

 $\sigma_{n,\kappa}$)²

Ø37)

one of the points. Using the strain ratio $\theta = \phi_2/\phi_1 (\phi_1 > \phi_2)$ [9], relation between equivalent and through-thickness strain is given as

$$\underline{2}^{\mathsf{r}} 1 + \theta + \theta^2$$

In modified hydraulic bulging with no draw in allowed, strain ratio $\boldsymbol{\theta}$ can only obtain values between 0 (plane strain) and 1 (equibiaxial stretching). Since the values

of function given by Eq \sqrt{uation} (17) are monotonically de-

3 ϕ_3 at $\theta = 0$ to the value ϕ_3 at

1

 $\phi_{eq} = 1 + \theta$

$$\begin{cases} F_{\kappa} \cos(\theta_{\kappa}) - & \theta = 1, \text{ i.e. equivalent strain} \\ \rho_{\kappa} R^2 \pi \mu [\cos(2\theta_{\kappa}) & 2 \end{cases}$$

$$\begin{array}{c} 1 & \phi_{eq} \in [1] \\ -2\sin(2\theta_{\kappa})] & .\sqrt{3} \\ \phi_{3} \end{array}$$

(18) present

In the present case of modified hydraulic bulging on the sphere two cases of strain were considered

Friction factor was taken to be $\mu = 0.2$. Division of meridional force with through-thickness area A_E at point *E* results in stress

 $\sigma_{m,E} =$

 $-\left[\theta_{E} - \frac{1}{2}\sin(2\theta_{E})\right]$

1. For the points *K* and *O* equibiaxial stretching $\beta = 1$ was assumed

2. For the point *E* plane strain $\theta = 0$ was assumed

(12)

$$\frac{F_E}{A_E} = \frac{F_E}{\pi s_E (2r_E + s_E r_E R)}$$

 $R^2\pi\{[\theta_\kappa$

)]}}

This assumption resulted in true stress-strain curve recorded for 1 mm thick Al 99.5, shown in Figure 4

 $\begin{array}{c} & & & \\ & & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & & \\ & & & \\ & & & \\ & & & \\ & & & & \\ & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & &$

Since point E lies on the sphere, radial and meridional curvature are the same and from equation (1) circular strain is calculated as

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Dogo Rangsang Research Journal ISSN : 2347-7180 $\sigma_{c,E} = \frac{p_K}{\sigma_{SE}} \sigma_{SE}$

Normal pressure at point E is the same as normal pressure at point K (see Section 3.1). So the equivalent stress at point E is given as

 $= \frac{1}{[(\sigma 4)]}$

σekv,E

$$+(\sigma_{c,E} - \sigma_{n,E})$$

m,E

-**σ**c,E

3.3 STRESS AT POINT O

Assumption of equal meridional and circular stress at the pole $\sigma_m = \sigma_c$ (commonly used for the stress at the pole in hydraulic bulging [3]), known thickness s_0 and meridional and circular radius equal to the radius of the sphere *R*; substituded into membrane equation (1) give in plane stress at point *O* as

where p_{K} is same contact pressure as in Section 3.1. Normal stress at point *O* equals normal stress at point *K* in Section 3.1 so equivalent stress at point *O* equals

$$\sigma_{ekv,O} = \left| \sigma_{m,c,O} - \sigma_{n,O} \right| \tag{16}$$

3.4 STRAIN IN MODIFIED BULGING

In order to record true stress-strain diagram, equivalent strain has to be determined for each one of the points K, E and O. To simplify the measurements, equivalent strain is determined from through-thickness strain at each

4.2 RATE OF DEFORMATION

Rate of deformation was calculated by dividing difference of strain, at point of the most intensive yielding *E*, for two subsequent pressures; 23 bar and 24 bar with respect to time needed for pressure to rise from 23 bar to 24 bar. Maximum rate of deformation of $\phi^{-} = 0.25 \ s^{-1}$ in modified hydraulic bulging lies on the lower boundary of the considered strain rates in the atlas of true stress - strain curves [5].

5 CONCLUSIONS

Modified hydraulic bulging on the sphere represents a new method of acquisition of true stress - strain curve. For the purpose of pre testing of the method, only one mate- rial; 1 mm thick aluminium Al 99.5 was used to record the curve shown in Figure 4.

5.1 ADVANTAGES

One great benefit of the proposed method is that; in contrast to classical hydraulic bulging, geometry of the edge UGC Care Journal Vol-10 Issue-06 No. 01 June 2020 (13)



Figure 4: Schematic representation of modified hydraulic bulging

4 EXPERIMENTAL SETUP

4.1 GEOMETRY OF THE TOOLING

In order to conduct modified hydraulic bulging on the

sphere, special tooling was designed and made. Bulging diameter was chosen to be D = 120 mm, radius of the

corner of the die r = 5 mm and radius of the sphere R = 20.6 mm as shown in Figure 5

Acquisition of the force and geometry was performed at high pressures; from 19 bar to the limiting pressure at 24 bar. This pressure span is uniformly covered by strain ranging from 0.1 to 0.7 as follows;

- at point O strains are from 0.1 to 0.2
- at point K strains are from 0.2 to 0.3
- at point *E* strains are from 0.3 to 0.7

of the die and die clamping diameter, both shown in Figure 5, are irrelevant for recorded true stress - strain curve. Moreover, future investigations will include stress and strain acquisition at all points from axis of symmetry to the point *B* shown in Figure 2.

Interesting detail to be considered is a possibility to perform modified hydraulic bulging without a force transducer, using only assumption of the membrane stress at point *B*. Since there is no bending stress, exerted force onto the surface of radius r_B by presure p, equals the force *F* on the sphere, shown in Figure 2.

For all of this purposes, a finite element model of modified hydraulic bulging on the sphere is being developed in order to support the future research.

5.2 DRAWBACKS

This method needs further development regarding; acquisition of strains larger than $\phi > 0.3$ (assumption on strain ratio θ), quantitative determination regarding influence of friction onto the resulting true stress - strain curve and comparison to the results of standard methods [5].

In this paper comparison of the true stress - strain diagram

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shown in Figure 4 and stress - strain diagram for Al 99.5 was not performed since only one material and one value of thickness was considered.

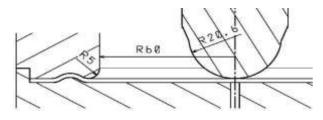


Figure 5: Central part of the tooling; modified hydraulic bulging on the sphere

Friction was only tested for the extremal values of $\mu = 0$ and $\mu = 0.58$ and found that it contributes very little to the *F*. stress. However, more detailed analysis is required.

Assumption on strain ratio $\beta = 0$ at point *E* holds only for the ultimate pressures. Since equivalent strain is additive value regarding strain path, further study of modified hydraulic bulging has to consider strain history at point

E. At this moment only the upper limit of cumulative equivalent strain at point E has been assumed yielding data shown in Figure 4.

Although modified hydraulic bulging gives more direct way to acquire stress, future work has to study strain distribution in greater detail relying on available photogrammetric methods [10].

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